# FROM THE YUKAWA TO THE EFIMOV ATTRACTION

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BARAMATER

#### THE YUKAWA POTENTIAL

 In the 1930s, Hideki Yukawa showed that the exchange of a massive boson field between two particles induces a Coulomb potential exponentially screened by the boson mass.





Mediated potential  $V(r) = -g^2 \frac{e^{-kmr}}{r}$ 

#### THE YUKAWA POTENTIAL

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- Coulombadetaatie Leftebacteranaanetiamige exphanagenetiamige exphanagenetiamige exphanaged particles •
- Tail of the nuclear force: exchange of mesons (135 MeV) between nuclear force: exchange of mesons ( $m \ge 135$  MeV) between nucleons •

#### THE EFIMOV POTENTIAL

 In the 1970s, Vitaly Efimov showed that a universal three-body force arises between three resonantlyinteracting particles.

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 $V(R) = -\frac{\hbar^2}{2M} \frac{1.006^2}{R^2}$ 

3 partitleseef onass M

mass

R

 $V(r) = -\frac{\hbar^2}{2m} \frac{0.414^2}{r^2}$ 

May also be viewed as a two-body mediated potential

Diagrammatic point of view



$$\frac{\frac{1}{g} + \sum_{k}^{\Lambda} \frac{1}{\frac{\hbar^2 k^2}{2\mu}} = \frac{2\mu}{4\pi\hbar^2} \frac{1}{a}}{<0} > 0$$

Scattening hength Reduced mass  $\mu = \left(\frac{1}{m} + \frac{1}{M}\right)^{-1}$ 

#### Diagrammatic point of view



$$\frac{1}{g} + \sum_{k}^{\Lambda} \frac{1}{\frac{\hbar^2 k^2}{2\mu}} = \frac{2\mu}{4\pi\hbar^2} \frac{1}{a}$$

$$< 0 \qquad > 0$$



Scatteringnength Reduced mass  $\mu = \left(\frac{1}{m} + \frac{1}{M}\right)^{-1}$ 





## SYSTEMS EXHIBITING THE EFIMOV ATTRACTION

#### arXiv:1610.09805

**IOP** Publishing

Reports on Progress in Physics

Rep. Prog. Phys. 00 (2017) 000000 (77pp)

Review

#### Efimov physics: a review

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#### Abstract

This article reviews theoretical and experimental advances in Efimov physics, an array of quantum few-body and many-body phenomena arising for particles interacting via short-range resonant interactions, that is based on the appearance of a scale-invariant three-body attraction theoretically discovered by Vitaly Efimov in 1970. This three-body effect was originally proposed to explain the binding of nuclei such as the triton and the Hoyle state of carbon-12, and later considered as a simple explanation for the existence of some halo nuclei. It was subsequently evidenced in trapped ultra-cold atomic clouds and in diffracted molecular beams of gaseous helium. These experiments revealed that the previously undetermined three-body parameter introduced in the Efimov theory to stabilise the three-body consequences of the Efimov attraction have been since investigated theoretically, and are expected to be observed in a broader spectrum of physical systems.

#### MOTIVATION

Connection between Yukawa and Efimov mediated potentials?

Yukawa

Efimov

Many-body Bosons are created/absorbed Three-body Boson always there



boson









# Bose polaron recently observed



Jørgensen et al, PRL 117, 055302 (2016)



Ming-Guang Hu et al, PRL 117, 055301 (2016)













The (Bogoliubov) quasi-particles excitations of the BEC can mediate a Yukawa interaction

excitatio

To second-order in perturbation

 $1/\xi$ 

Theory:  
$$V(r) \propto -g^2 n_0 \frac{e^{-\sqrt{2r}}}{r}$$

BEC coherence length  $\xi = \frac{1}{\sqrt{8\pi n_0 a_B}}$ 

Ex: Helium-3 impurities in Helium-4 J. Bardeen, G. Baym, and D. Pines, Phys Rev 156, 207 (1067)



The (Bogoliubov) quasi-particles excitations of the BEC can also mediate an Efimov interaction

excitatio

g

#### Non-perturbative

#### **NON-PERTURBATIVE METHOD:** TRUNCATED BASIS



on

Bogoliubov $b_0 = \sqrt{N_0}$ condensateappproximati $b_k = u_k \beta_k - v_k \beta_k^{\dagger}$ Bogoliu Bogoliubov excitation

$$H = \begin{bmatrix} E_{0} + \sum_{k} E_{k} \beta_{k}^{\dagger} \beta_{k} \\ + \sum_{k} (\varepsilon_{k} + gn_{0})c_{k}^{\dagger}c_{k} \\ + \sqrt{N_{0}} \frac{g}{V} \sum_{k,p} (u_{p} \beta_{-p}^{\dagger} - v_{p} \beta_{p})c_{k+p}^{\dagger}c_{k} + h.c. \\ Fröhlich \\ + \frac{g}{V} \sum_{k,k',p} (u_{k'-p} u_{k'} \beta_{k'-p}^{\dagger} \beta_{k'} + v_{k'-p} v_{k'} \beta_{p-k'} \beta_{-k'}^{\dagger})c_{k+p}^{\dagger}c_{k} \\ + \frac{g}{V} \sum_{k,k',p} (u_{k'-p} v_{k'} \beta_{k'-p}^{\dagger} \beta_{-k'}^{\dagger} + v_{k'-p} u_{k'} \beta_{p-k'} \beta_{-k'})c_{k+p}^{\dagger}c_{k} \\ + \frac{g}{V} \sum_{k,k',p} (u_{k'-p} v_{k'} \beta_{k'-p}^{\dagger} \beta_{-k'}^{\dagger} + v_{k'-p} u_{k'} \beta_{p-k'} \beta_{k'})c_{k+p}^{\dagger}c_{k} \\ + \frac{g}{V} \sum_{k,k',p} (u_{k'-p} v_{k'} \beta_{k'-p}^{\dagger} \beta_{-k'}^{\dagger} + v_{k'-p} u_{k'} \beta_{p-k'} \beta_{k'})c_{k+p}^{\dagger}c_{k} \\ + \frac{g}{V} \sum_{k,k',p} (u_{k'-p} v_{k'} \beta_{k'-p}^{\dagger} \beta_{-k'}^{\dagger} + v_{k'-p} u_{k'} \beta_{p-k'} \beta_{k'})c_{k+p}^{\dagger}c_{k} \\ + \frac{g}{V} \sum_{k,k',p} (u_{k'-p} v_{k'} \beta_{k'-p}^{\dagger} \beta_{-k'}^{\dagger} + v_{k'-p} u_{k'} \beta_{p-k'} \beta_{k'})c_{k+p}^{\dagger}c_{k} \\ + \frac{g}{V} \sum_{k,k',p} (u_{k'-p} v_{k'} \beta_{k'-p}^{\dagger} \beta_{-k'}^{\dagger} + v_{k'-p} u_{k'} \beta_{p-k'} \beta_{k'})c_{k+p}^{\dagger}c_{k} \\ + \frac{g}{V} \sum_{k,k',p} (u_{k'-p} v_{k'} \beta_{k'-p}^{\dagger} \beta_{-k'}^{\dagger} + v_{k'-p} u_{k'} \beta_{p-k'} \beta_{k'})c_{k+p}^{\dagger}c_{k} \\ + \frac{g}{V} \sum_{k,k',p} (u_{k'-p} v_{k'} \beta_{k'-p}^{\dagger} \beta_{-k'}^{\dagger} + v_{k'-p} u_{k'} \beta_{p-k'} \beta_{k'})c_{k+p}^{\dagger}c_{k} \\ + \frac{g}{V} \sum_{k,k',p} (u_{k'-p} v_{k'} \beta_{-k'} \beta_{-k'} \beta_{-k'} + v_{k'-p} u_{k'} \beta_{p-k'} \beta_{k'})c_{k+p}^{\dagger}c_{k} \\ + \frac{g}{V} \sum_{k,k',p} (u_{k'-p} v_{k'} \beta_{-k'} \beta_{-k'} + v_{k'-p} u_{k'} \beta_{-k'} \beta_{-k'}$$

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#### **NON-PERTURBATIVE METHOD: TRUNCATED BASIS**



Impurity creation operator

Excitatio n creation operator

BEC groun d state

#### EQUATIONS

Coupled equiations  $d_{pand}$  and  $d_{pq'}$ 

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 $F_q = g \frac{1}{V} \sum_{k} u_k \alpha_{q,k-q}$ 

Three-body Integral equation  $\frac{1}{T_q(E)}F_q + \frac{1}{V}\sum_k \frac{u_k^2 F_{k-q}}{E_k + \varepsilon_{|k-q|} + \varepsilon_q - E} = \frac{2n_0}{2\varepsilon_q - E}F_q$ 

$$-\frac{1}{T_q(E)} = \frac{2\mu}{4\pi\hbar^2} \frac{1}{a} + \frac{1}{V} \sum_k \left( \frac{u_k^2}{E_k + \varepsilon_{|k-q|} + \varepsilon_q - E} - \frac{1}{\varepsilon_k + \varepsilon_k} \right)$$

#### **RESULT: POLARONIC POTENTIAL**

Effective potential (Born-Oppenheimer) between polarons:

$$V(r) = \frac{\hbar^2 \kappa^2}{2\mu} \qquad \qquad \frac{1}{a} - \kappa + \frac{1}{r} e^{-\kappa r} + \frac{8\pi n_0}{\kappa^2} = 0$$

For small  $k \leq 0$ V(r) V(r) V(r) V(r)V(r)

#### At resonance $\pm \infty$

 $(a_B \rightarrow 0)$ 













# POSSIBLE EXPERIMENTAL OBSERVATIONS

Heavy impurities (e.g. <sup>133</sup>Cs) in a condensate of light bosons (e.g. <sup>7</sup>Li)

- Polaron RF spectroscopy, mean-field side to the polaron interaction
- Loss by recombination peak with the light bu

h<mark>ift of the loss</mark> i density

# CONCLUSIONS

arXiv:1607.04507

- The interaction between two impurities in a Bose-Einstein goes from a Yukawa attraction to an **Efimov attraction**.
- This attraction can bind the two impurities into one or several bipolarons, that asymptote to Efimov trimers of two impurities and a boson.
- The crossover region where bipolarons appear constitutes is an interesting few-to-many-body problem.