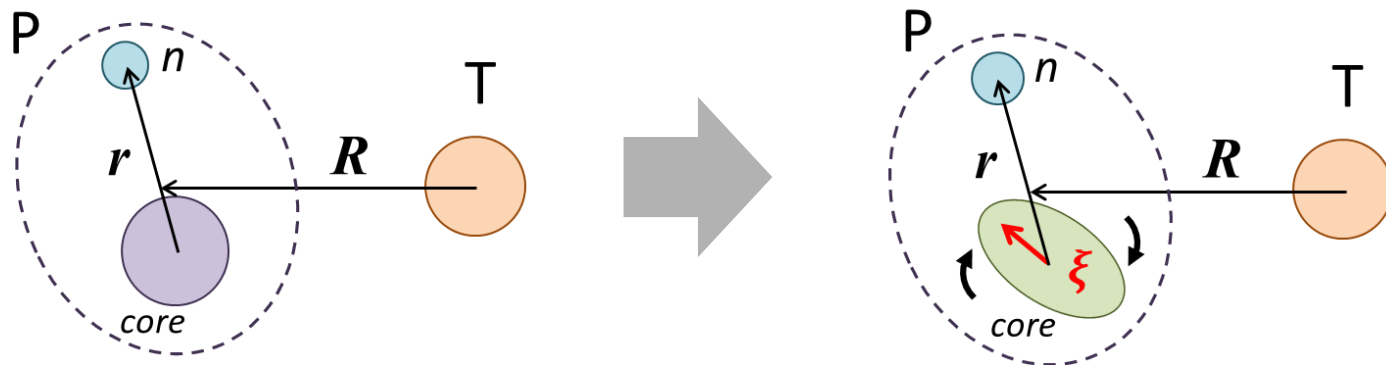


Dynamic and static core excitation effects on deformed halo nuclei

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¹RIKEN, ²RCNP, Osaka University, ³Kyushu University



10/March/2017

International Workshop on Quantum Many-Body Problems
in Particle, Nuclear, and Atomic Physics
Duy Tan University, Danang City, Vietnam

Discovery of halo

VOLUME 55, NUMBER 24

PHYSICAL REVIEW LETTERS

9 DECEMBER 1985

Measurements of Interaction Cross Sections and Nuclear Radii in the Light p -Shell Region

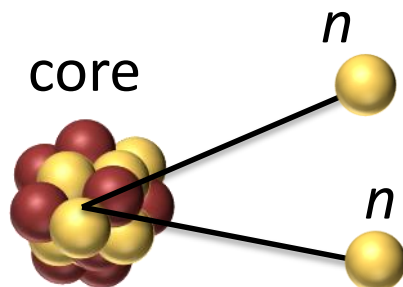
I. Tanihata,^(a) H. Hamagaki, O. Hashimoto, Y. Shida, and N. Yoshikawa
Institute for Nuclear Study, University of Tokyo, Tanashi, Tokyo 188, Japan

K. Sugimoto,^(b) O. Yamakawa, and T. Kobayashi
Nuclear Science Division, Lawrence Berkeley Laboratory, University of California, Berkeley, California 94720

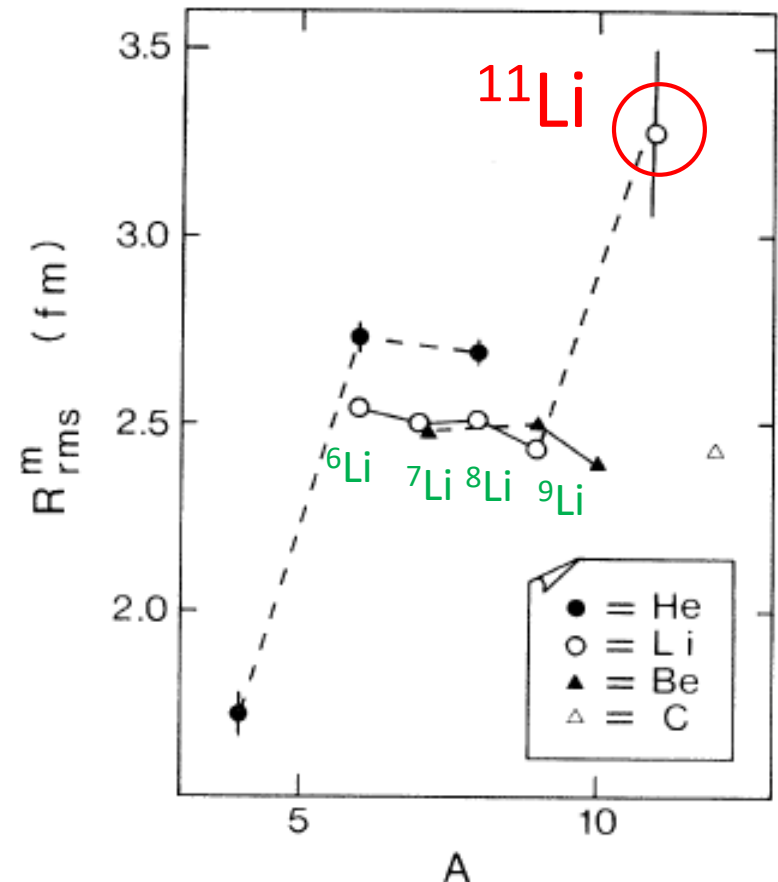
and
N. Takahashi

Neutron halo

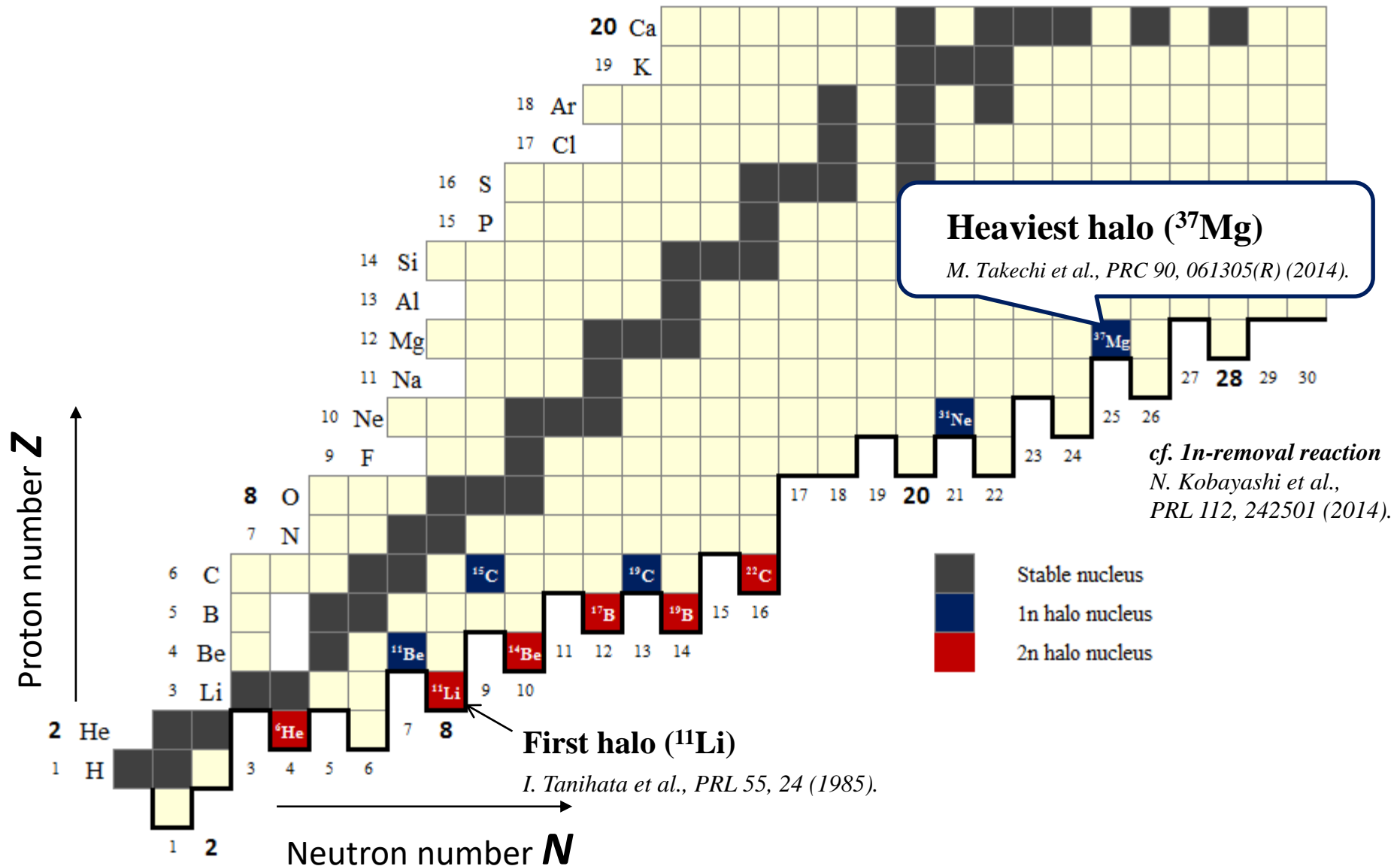
One or two neutron(s) surround very far from a core nucleus.



A large cross section was measured in Li isotopes.



Recent development of reaction cross sections (σ_R)



Systematic analysis of σ_R

Experiment

Total reaction cross sections (σ_R) were measured systematically.

M. Takechi et al., PRC 90, 061305(R) (2014).

M. Takechi et al., PLB 707, 357 (2012).

Theory

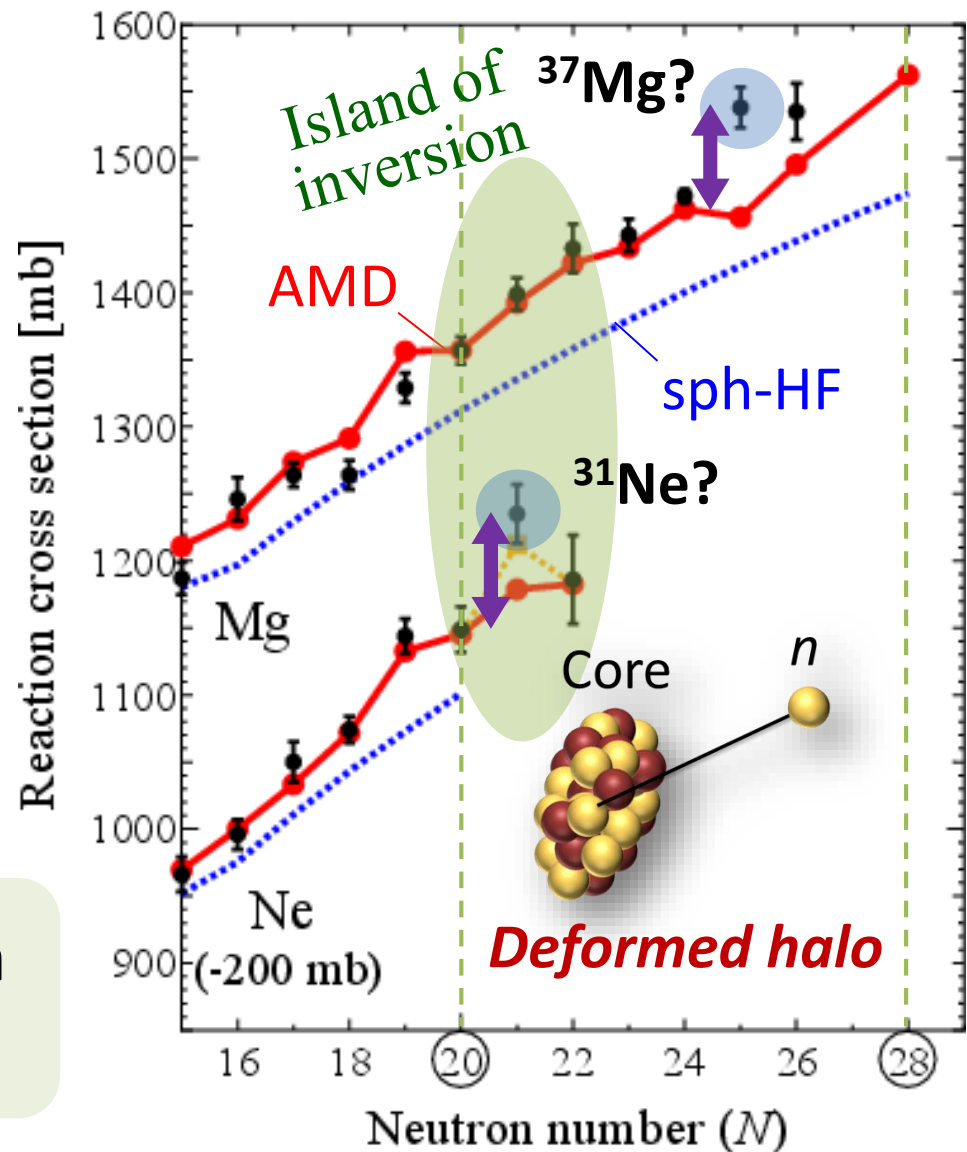
The σ_R were analyzed in the microscopic framework based on AMD and the double folding model.

S. Watanabe et al., PRC 89, 044610 (2014).

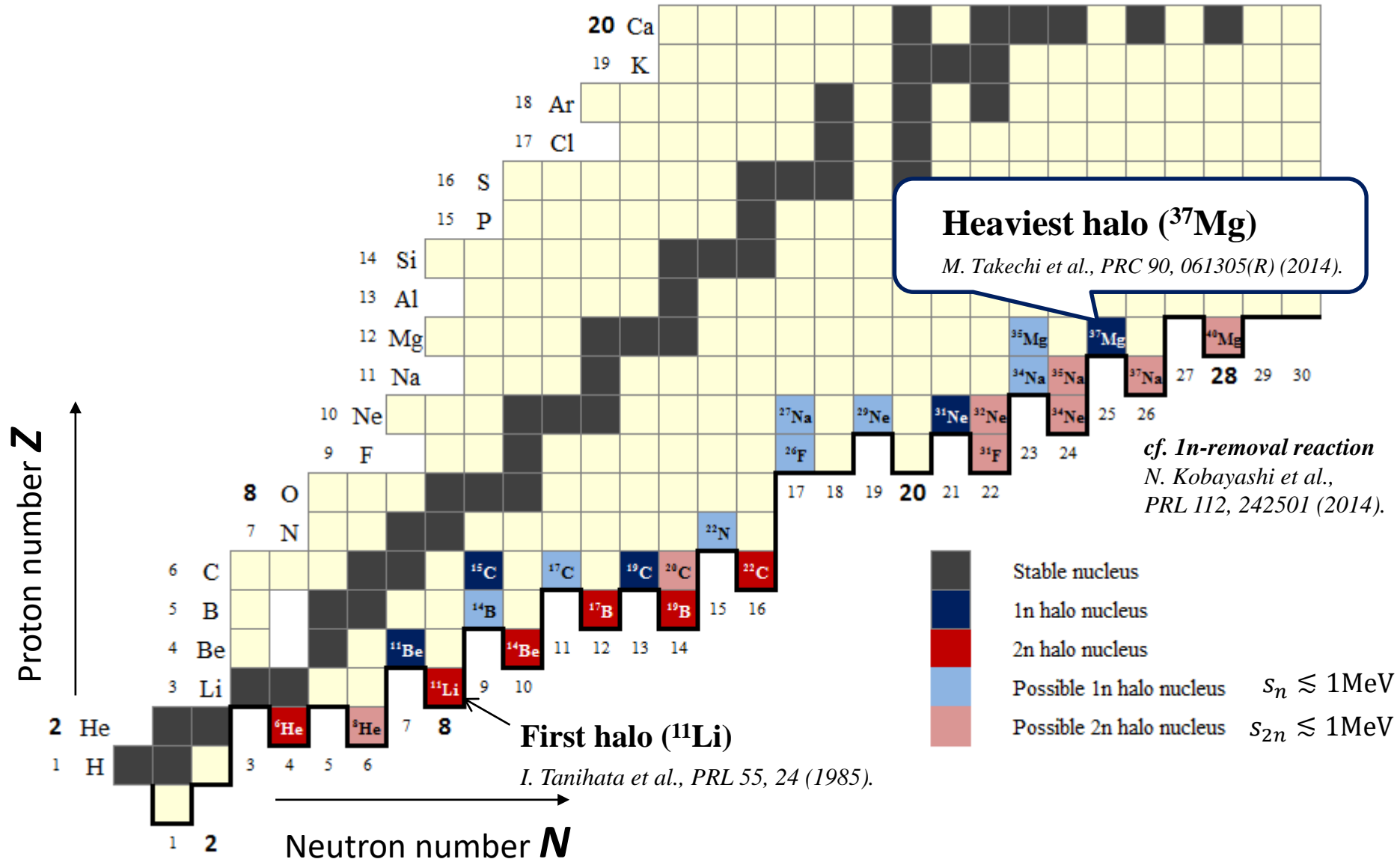
K. Minomo et al., PRL 108, 052503 (2012).

$19 \leq N$: Largely deformation

^{31}Ne , ^{37}Mg : **Deformed halo**



Recent development of reaction cross sections (σ_R)



Recent development of reaction cross sections (σ_R)

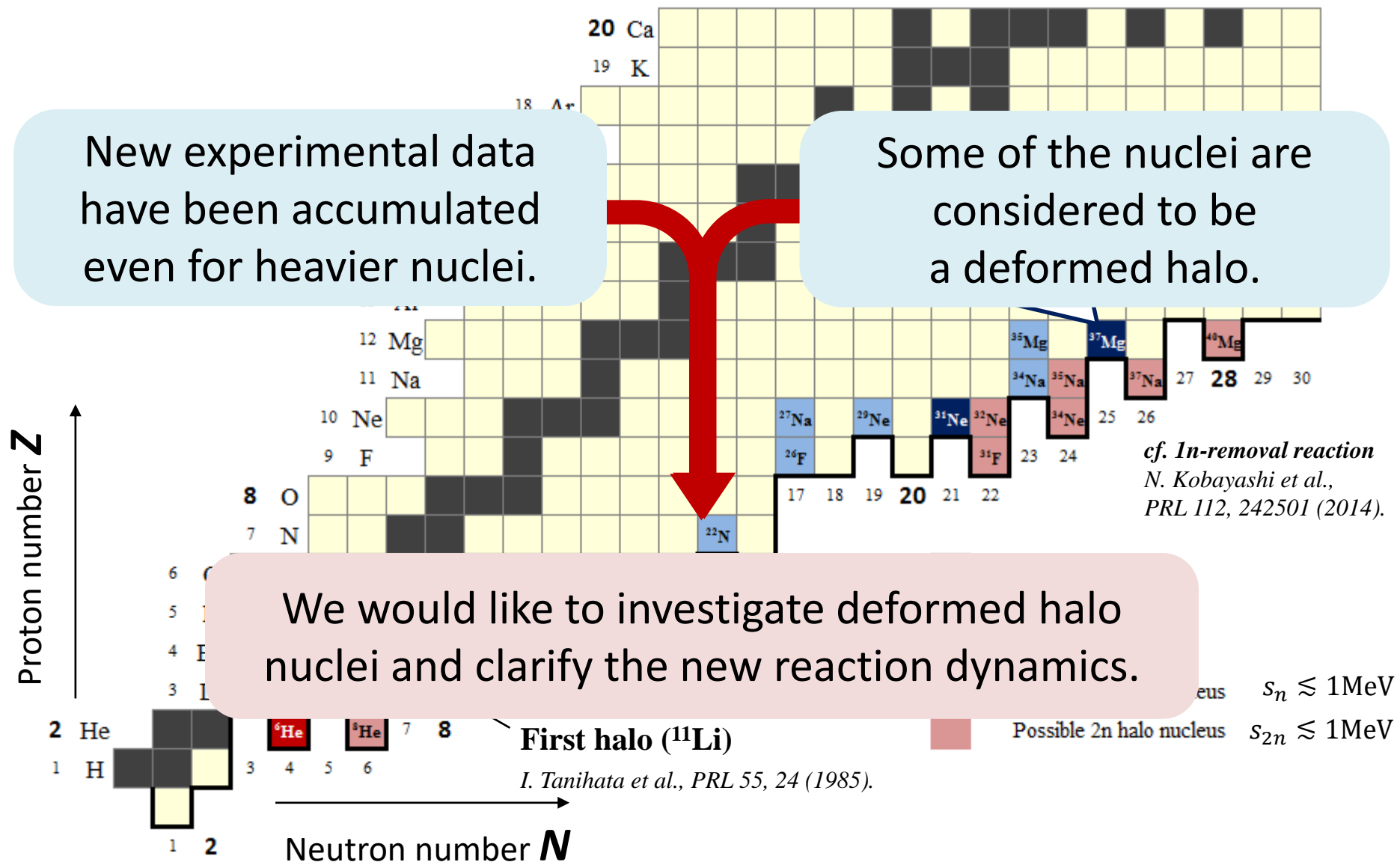


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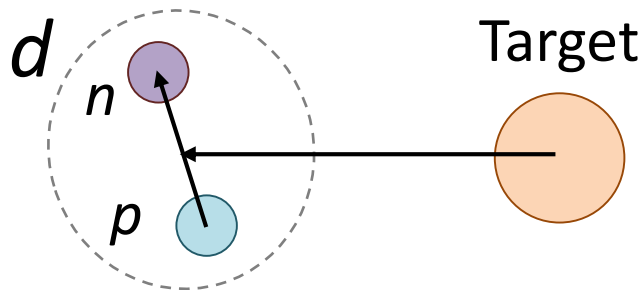
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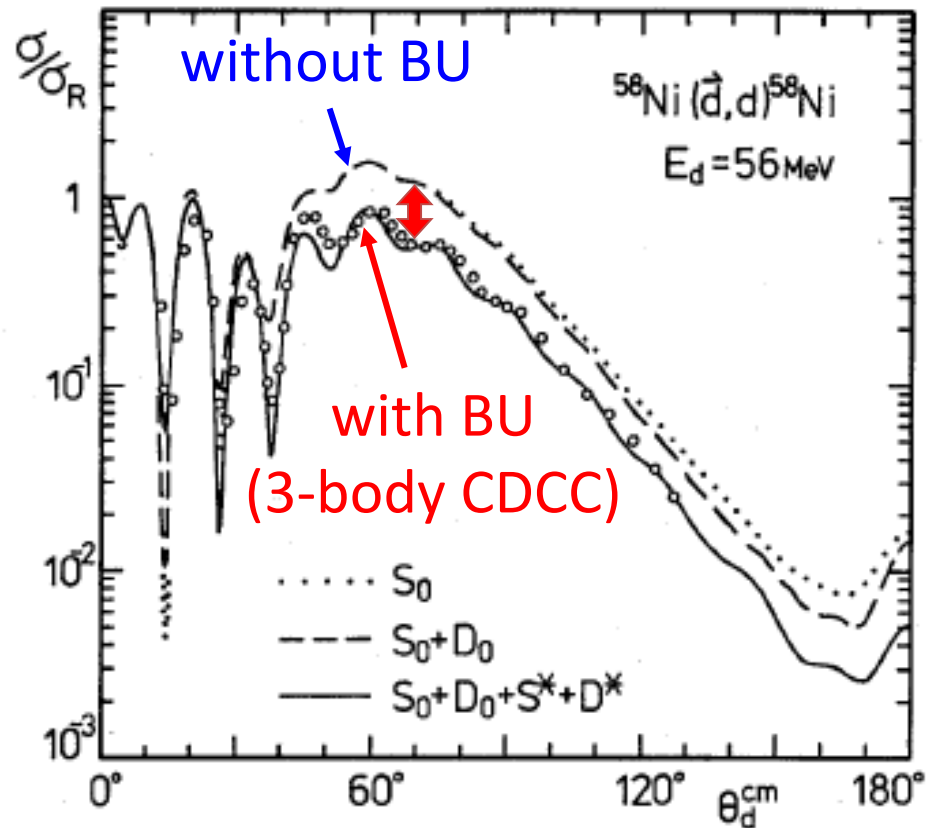
CDCC and Breakup effects

CDCC Continuum Discretized Coupled Channels

- ✓ CDCC is a fully quantum mechanical method for treating BU effects.
- ✓ CDCC was born as a theory for d -scattering



CDCC has been widely applied to many kinds of three-body scattering. (Ex: core + n + T)



First application of 3-body CDCC

M. Yahiro, Y. Iseri, H. Kameyama, M. Kamimura, and M. Kawai, Prog. Theor. Phys. Suppl. No. 89 (1986), 32.

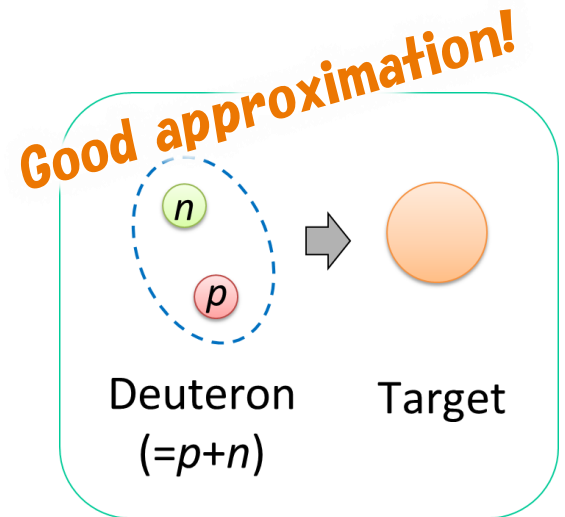
General few-body approaches (CDCC, Faddeev, DWBA, ...)

A core nucleus is assumed to be inert.

😊 **Good approximation** for d scattering

- Neutron and proton cannot get excited in the energy scale of our interest.

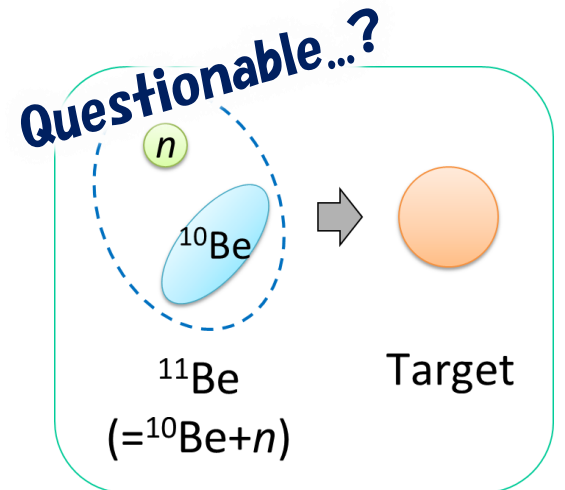
cf. ^6He scattering (α core is *inert*)



☹ **Questionable** for heavier systems

- Different core and valence states are coupled with each other.

ex. ^{11}Be , ^{37}Mg (Core is *deformed*)



Immediate work: To develop CDCC for treating core excitation.

Important DoF: Core excitation

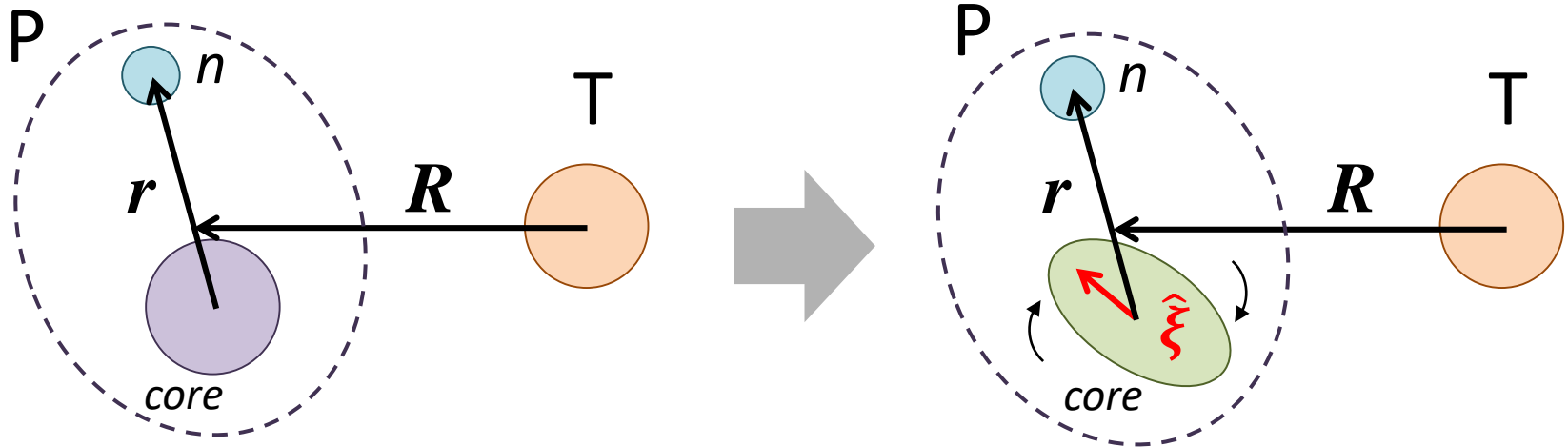
We will introduce the core coordinate ($\hat{\xi}$).

Standard 3-body CDCC

$$\Psi = \Psi(R, r)$$

3-body CDCC with core Ex.

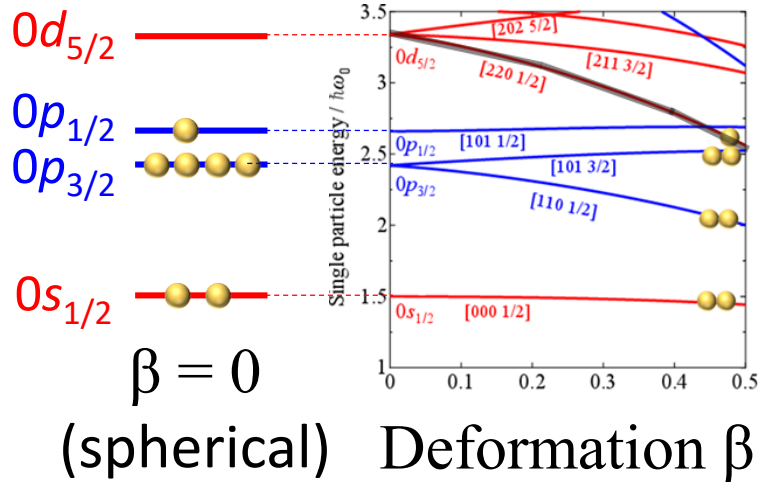
$$\Psi = \Psi(R, r, \hat{\xi})$$



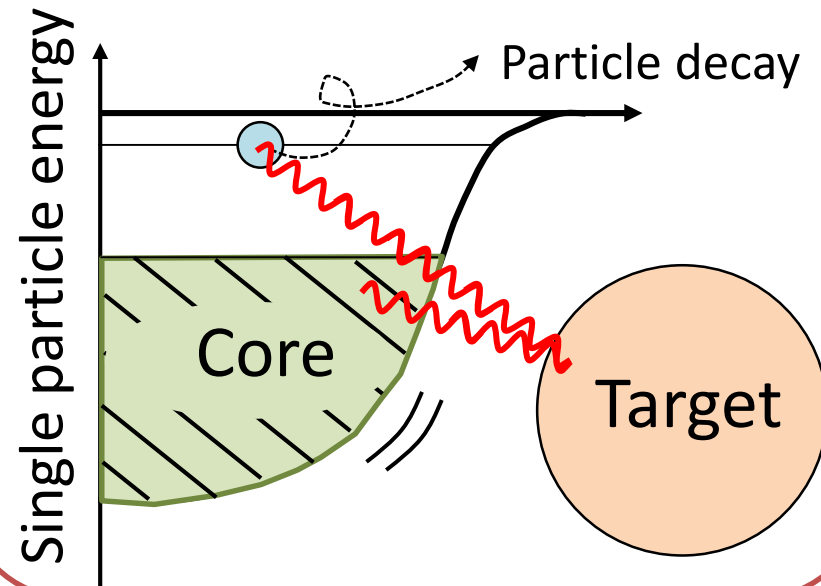
What kind of physics appears
by taking into account the **core excitation**?

Static and Dynamic Core Excitation Effects

Static core excitation



Dynamic core excitation



- Single particle energy changes
- Coupled to several core states

- BU due to neutron excitation
- BU due to core excitation

Competition

I would like to understand “Static” and “Dynamic” core excitation effects simultaneously.

Purpose

Final goal

To develop the **CDCC** method for treating core-excitation effects explicitly.

- Investigation of scattering of deformed halos
- Application for cluster physics

DWBA: First order approximation of CDCC

Present goal

To develop
DWBA (Distorted Wave Born Approximation)
for treating core excitation.

- Good prototype (DWBA = CDCC for $\hat{V} \sim 0$)
- Simple estimation of core-excitation effects in the BU reaction

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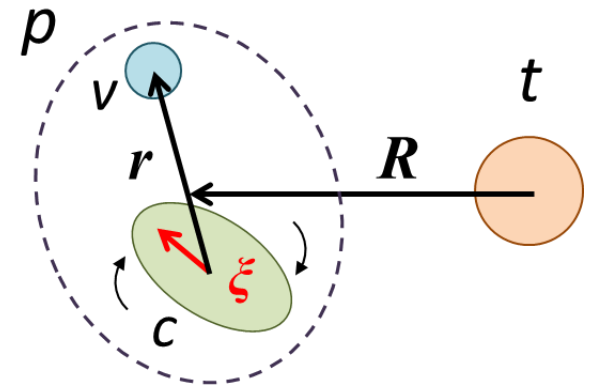
Total Hamiltonian

We explicitly introduce the core DoF ($\hat{\xi}$).

$$H_{\text{tot}} = K_{\mathbf{R}} + V_{vt}(R_{vt}) + V_{ct}(\mathbf{R}_{ct}, \hat{\xi}) + h_p$$

$$h_p = K_r + V(\mathbf{r}, \hat{\xi}) + h_c(\hat{\xi})$$

→ I will show you the actual interactions later.



First, we should understand projectile part (h_p).

→ Let's take ^{11}Be as an example.

Particle Rotor Model (PRM)

✓ ^{11}Be is described as
 n (Particle) + active core (Rotor)

■ Hamiltonian

$$H = -\frac{\hbar^2}{2\mu} \nabla_{\mathbf{r}}^2 + V(\mathbf{r}, \hat{\xi}) + \frac{\hbar^2 I_c^2}{2J}$$

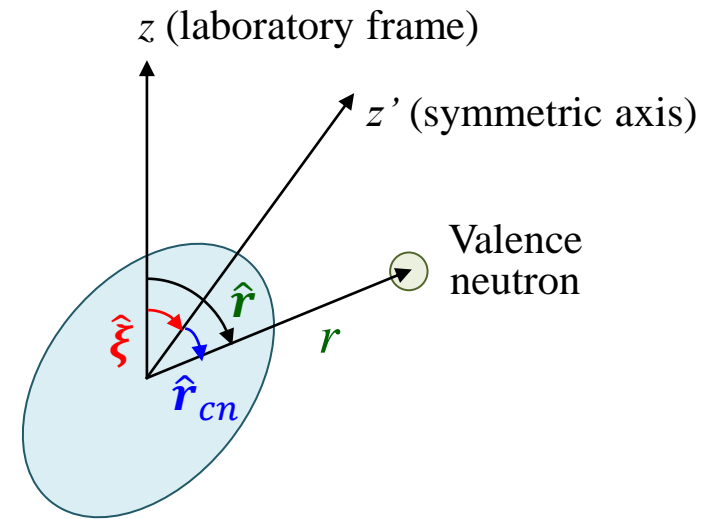
$$\approx \underbrace{-\frac{\hbar^2}{2\mu} \nabla_{\mathbf{r}}^2 + V_0(r)}_{\text{Spherical Hamiltonian}} + \underbrace{-R_0 \beta_2 \frac{dV_0(r)}{dr} Y_{20}(\hat{\mathbf{r}}_{cn})}_{\text{Deformation term}} + \underbrace{\frac{\hbar^2 I_c^2}{2J}}_{\text{Internal Hamiltonian of core}}$$

Spherical Hamiltonian
 $\equiv H_{\text{sph}}$

Deformation term
 $\equiv V_{\text{def}}$

Internal Hamiltonian of core
 $\equiv H_{\text{rot}}$

$$H \equiv H_{\text{sph}}(\mathbf{r}) + V_{\text{def}}(\hat{\mathbf{r}}_{cn}) + H_{\text{rot}}(\hat{\xi})$$



How to solve this problem?

$$[H - E]\Psi_{JM}(\mathbf{r}, \hat{\xi}) = 0$$

$$H = H_{\text{sph}}(\mathbf{r}) + V_{\text{def}}(\mathbf{r}, \hat{\xi}) + H_{\text{rot}}(\hat{\xi})$$

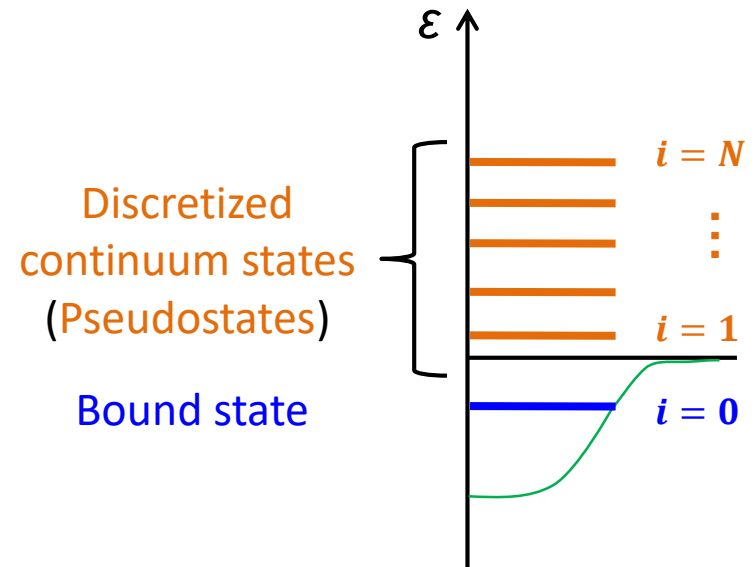
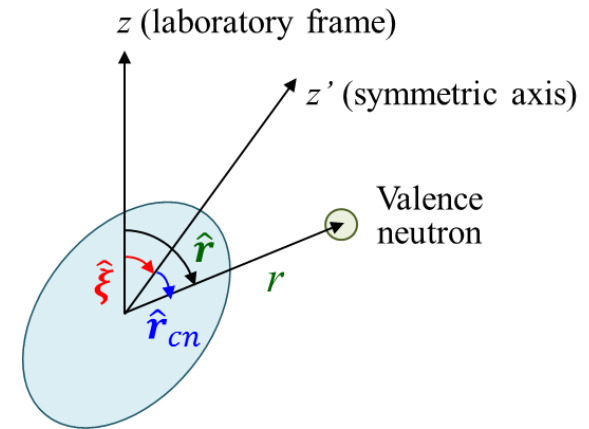
$$\Psi_{JM}(\mathbf{r}, \hat{\xi}) = \sum_{n\ell j} \sum_I \alpha_{n\ell j I} \left[\phi_{n\ell j}^{(\text{sph})}(\mathbf{r}) \otimes \Phi_I^{(\text{rot})}(\hat{\xi}) \right]_{JM}$$

$$\phi_{n\ell j}^{(\text{sph})}(\mathbf{r}) = \sum_{k=1}^N \beta_k \boxed{\varphi_k(r)} \psi_{(\ell s)j}(\hat{\mathbf{r}})$$

Gaussian basis

*E. Hiyama, Y. kino, M. Kamimura,
Prog. Part. Nucl. Phys. 51, 223 (2003).*

Diagonalization



Model Setting

Hamiltonian

$$H = -\frac{\hbar^2}{2\mu} \nabla_{\mathbf{r}}^2 + V_0(r) + V_{\text{def}}(\hat{\mathbf{r}}_{cn}) + \frac{\hbar^2 I_c^2}{2\mathcal{J}}$$

Parameter set of WS potential

$$V_{\text{WS}} = -54.239 \text{ MeV}, V_{\text{SO}} = -8.50 \text{ MeV}$$

$$R = 2.483 \text{ fm}, a = 0.65 \text{ fm}$$

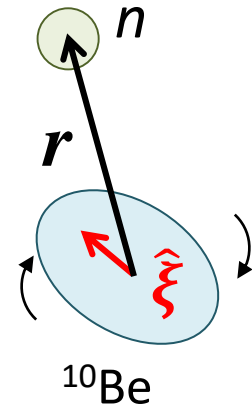
F.M. Nunes et al., NPA609 43 (1996).

^{10}Be core

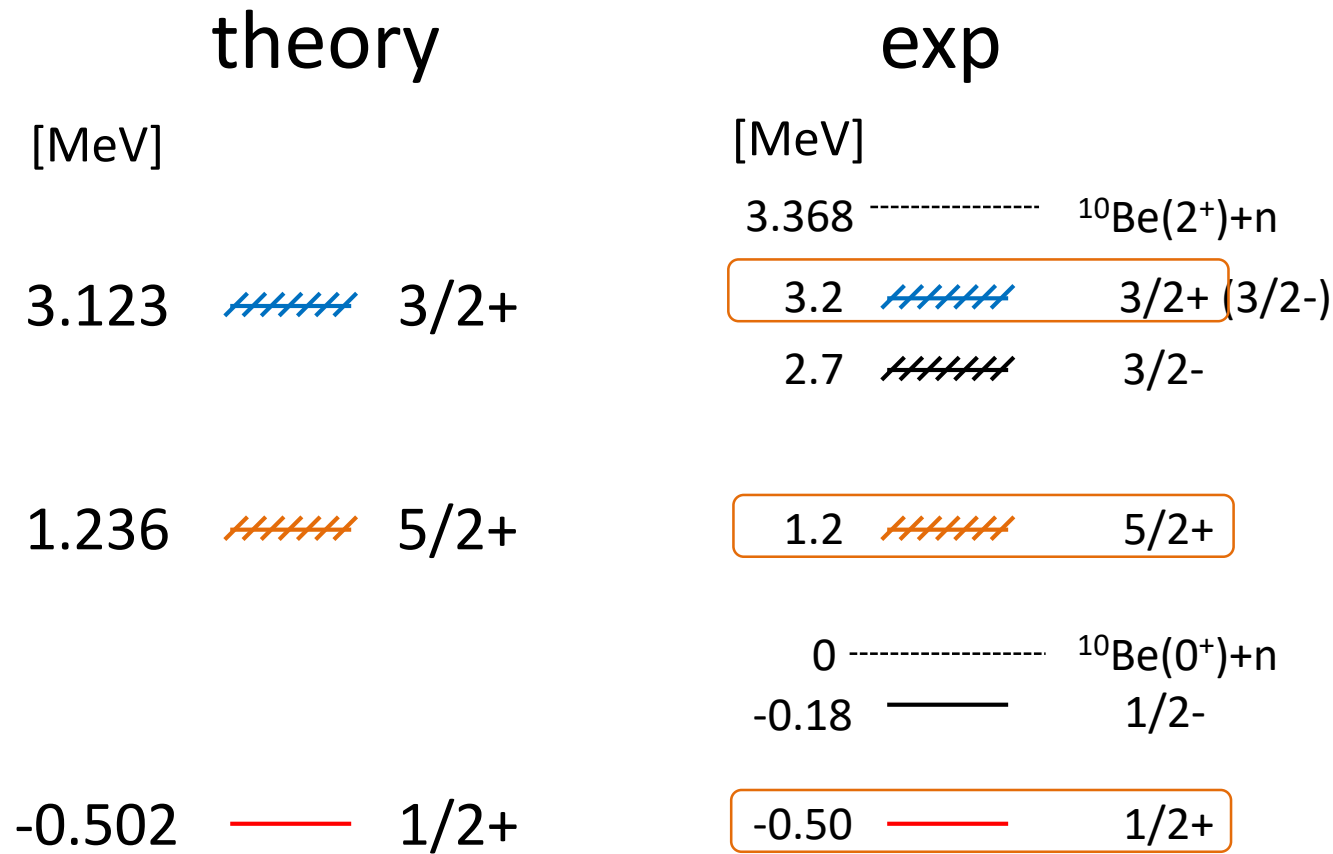
$$\beta_2 = 0.67, E(2+) = 3.368 \text{ MeV}$$

Model space

$$\ell = 0, 2 \quad I = 0, 2$$



Result of the positive-state energies



We will consider the $5/2^+$ and $3/2^+$ BU reactions.

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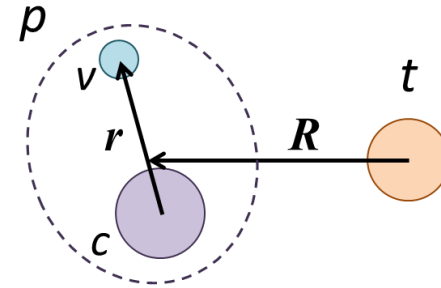
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DWBA (Distorted Wave Born Approximation)

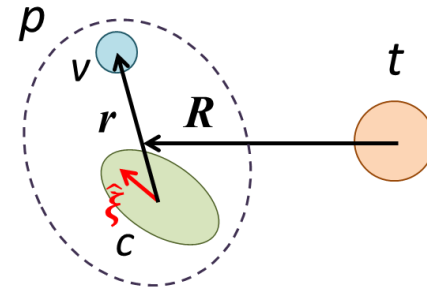
Standard DWBA (since 1950s)



$$T_{pt}^{J'M',JM}(\mathbf{K}', \mathbf{K}) = \left\langle \chi_{\mathbf{K}'}^{(-)}(\mathbf{R}) \Psi_{J'M'}^f(\mathbf{r}) \left| V_{vt}(R_{vt}) + V_{ct}(R_{ct}) \right| \chi_{\mathbf{K}}^{(+)}(\mathbf{R}) \Psi_{JM}^i(\mathbf{r}) \right\rangle$$

Extended DWBA

A. Moro and R. Crespo., PRC 85, 054613 (2012)



$$T_{pt}^{J'M',JM}(\mathbf{K}', \mathbf{K}) = \left\langle \chi_{\mathbf{K}'}^{(-)}(\mathbf{R}) \Psi_{J'M'}^f(\mathbf{r}, \hat{\xi}) \left| V_{vt}(R_{vt}) + V_{ct}(\mathbf{R}_{ct}, \hat{\xi}) \right| \chi_{\mathbf{K}}^{(+)}(\mathbf{R}) \Psi_{JM}^i(\mathbf{r}, \hat{\xi}) \right\rangle$$

static core excitation
dynamic core excitation

DWBA description of BU reaction

Extended DWBA

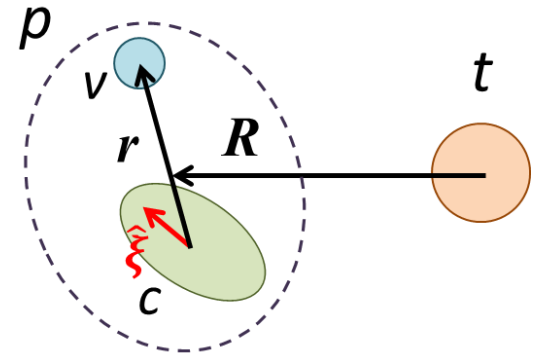
$$T_{pt}^{J'M',JM} = \left\langle \chi_{K'}^{(-)}(\mathbf{R}) \Psi_{J'M'}^f(\mathbf{r}, \hat{\xi}) \left| V_{vt}(R_{vt}) + V_{ct}(\mathbf{R}_{ct}, \hat{\xi}) \right| \chi_K^{(+)}(\mathbf{R}) \Psi_{JM}^i(\mathbf{r}, \hat{\xi}) \right\rangle$$

Multipole expansion of V_{ct}

$$V_{vt}(R_{vt}) + V_{ct}(\mathbf{R}_{ct}, \hat{\xi})$$

$$= V_{vt}(R_{vt}) + \sum_{\mathcal{LM}} V_{ct}^{(\mathcal{L})}(R_{ct}) Y_{\mathcal{LM}}^*(\hat{\mathbf{R}}_{ct}) Y_{\mathcal{LM}}(\hat{\xi})$$

$$= V_{vt}(R_{vt}) + V_{ct}^{(0)}(R_{ct}) + \sum_{\mathcal{L} > 0, \mathcal{M}} V_{ct}^{(\mathcal{L})}(R_{ct}) Y_{\mathcal{LM}}^*(\hat{\mathbf{R}}_{ct}) Y_{\mathcal{LM}}(\hat{\xi})$$

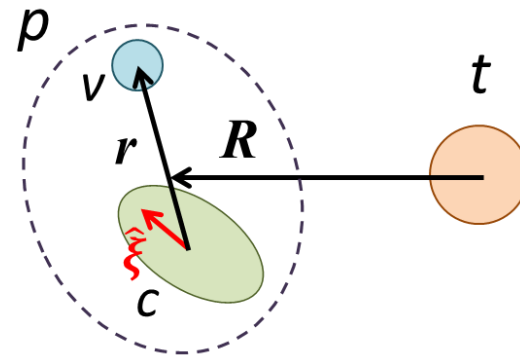


DWBA description of BU reaction

Extended DWBA

$$T_{pt}^{J'M',JM} = \left\langle \chi_{K'}^{(-)}(\mathbf{R}) \Psi_{J'M'}^f(\mathbf{r}, \hat{\xi}) \left| V_{vt}(R_{vt}) + V_{ct}(\mathbf{R}_{ct}, \hat{\xi}) \right| \chi_K^{(+)}(\mathbf{R}) \Psi_{JM}^i(\mathbf{r}, \hat{\xi}) \right\rangle$$

$$T_{pt}^{J'M',JM} = T_{\text{val}}^{J'M',JM} + T_{\text{corex}}^{J'M',JM}$$



Valence excitation

$$T_{\text{val}}^{J'M',JM} = \left\langle \chi_{K'}^{(-)}(\mathbf{R}) \Psi_{J'M'}^f(\mathbf{r}, \hat{\xi}) \left| V_{vt}(R_{vt}) + V_{ct}^{(0)}(R_{ct}) \right| \chi_K^{(+)}(\mathbf{R}) \Psi_{JM}^i(\mathbf{r}, \hat{\xi}) \right\rangle$$

→ Excite valence coordinate (\mathbf{r})

Core excitation

$$T_{\text{corex}}^{J'M',JM} = \left\langle \chi_{K'}^{(-)}(\mathbf{R}) \Psi_{J'M'}^f(\mathbf{r}, \hat{\xi}) \left| \sum V_{ct}^{(\mathcal{L})}(R_{ct}) Y_{\mathcal{LM}}^*(\hat{\mathbf{R}}_{ct}) Y_{\mathcal{LM}}(\hat{\xi}) \right| \chi_K^{(+)}(\mathbf{R}) \Psi_{JM}^i(\mathbf{r}, \hat{\xi}) \right\rangle$$

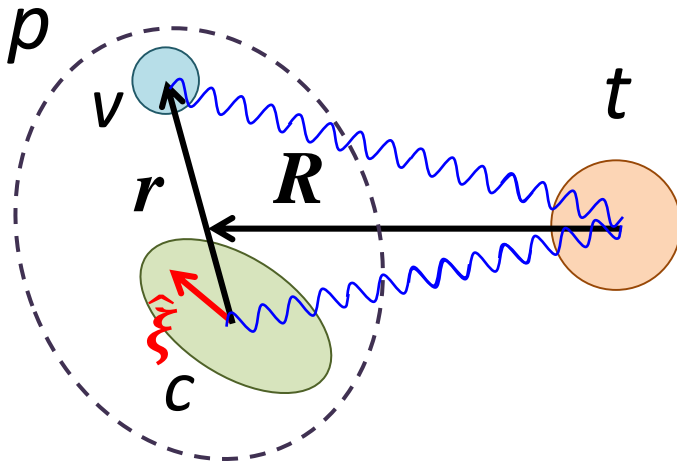
→ Excite core coordinate ($\hat{\xi}$)

Brief summary: Two types of BU mechanisms

$$T_{pt}^{J'M',JM}(K',K) = T_{\text{val}}^{J'M',JM}(K',K) + T_{\text{corex}}^{J'M',JM}(K',K)$$

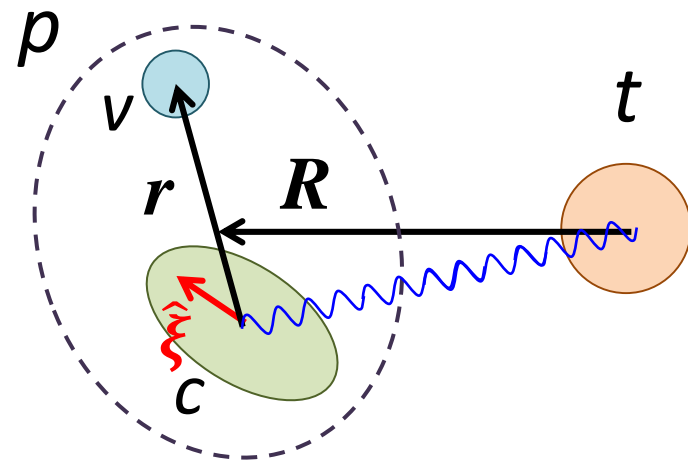
Valence excitation

interaction: $V_{vt}(R_{vt}) + V_{ct}^{(0)}(R_{ct})$



Core excitation

interaction: $\sum V_{ct}^{(\mathcal{L})}(R_{ct})Y_{\mathcal{LM}}^*(\hat{\mathbf{R}}_{ct})Y_{\mathcal{LM}}(\hat{\mathbf{r}})$



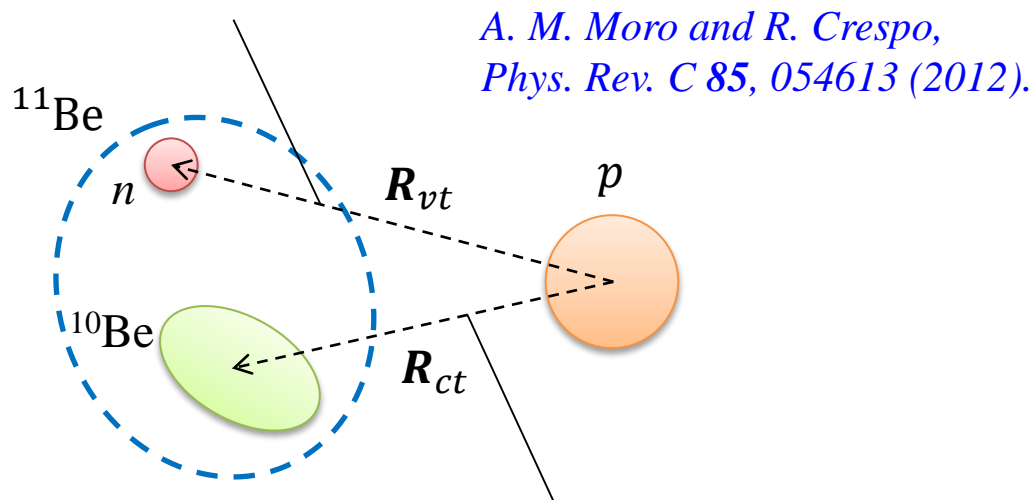
These two mechanisms compete with each other.

Model setting of reaction part

$^{11}\text{Be}+p$ at 63.7 MeV/nuc.

$$V_{nt}: V(r) = -45e^{-(r/1.484)^2}$$

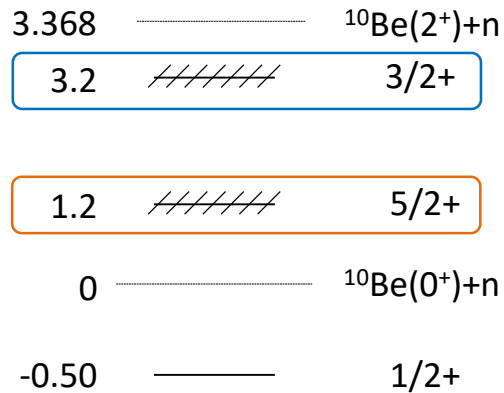
Determined to reproduce the realistic Faddeev calculation. (with CD Bonn)



V_{ct} : Phenomenological optical potential

B. A. Watson et al., Phys. Rev. 182, 977 (1969).

Resonant breakup cross section



$$\frac{d\sigma}{d\Omega} \propto \sum_{MM'} |T_{pt}^{J'M',JM}|^2$$

$$T_{pt}^{J'M',JM} = T_{\text{val}}^{J'M',JM} + T_{\text{corex}}^{J'M',JM}$$

Sorry.

Summary

We are developing our reaction model (CDCC) for the explicit treatment of core excitation.

First, we developed DWBA for treating core excitation.

Until now: Core is assumed to be **inert**.

From now: **Core excitation** will be a **key mechanism** in nuclear reaction.

➤ This effect appears in $^{11}\text{Be}+p$ scattering.

Future plan: Develop DWBA into CDCC.

Analyze ^{31}Ne , ^{37}Mg etc. (deformed halo)

^{15}C (spherical halo)

⇒ General properties of core-induced BU reaction?